

Generalized Darboux transformations for the KP equation with self-consistent sources

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Abstract

The KP equation with self-consistent sources (KPESCS) is treated in the framework of the constrained KP equation. This offers a natural way to obtain the Lax representation for the KPESCS. Based on the conjugate Lax pairs, we construct the generalized binary Darboux transformation with arbitrary functions in time t for the KPESCS which, in contrast with the binary Darboux transformation of the KP equation, provides a non-auto-Bäcklund transformation between two KPESCSs with different degrees. The formula for N-times repeated generalized binary Darboux transformation is proposed and enables us to find the N-soliton solution and lump solution as well as some other solutions of the KPESCS.

Keywords: KP equation with self-consistent sources(KPESCS), conjugate Lax pairs, Darboux transformation(DT), N-soliton solution, lump solution

1 Introduction

Soliton equations with self-consistent sources (SESCSs) are important models in many fields of physics, such as hydrodynamics, solid state physics, plasma physics, etc. [1-7]. Until now, much development has been made in the study of SESC. For example, in (1+1)-dimensional case, some SESC. such as the KdV, modified KdV, nonlinear Schrodinger, AKNS and Kaup-Newell hierarchies with self-consistent sources were solved by the inverse scattering method [1,2,3,6-10]. Recently, generalized binary Darboux transformations with

arbitrary functions in time t for some (1+1)-dimensional SESCOs, which offer a non-auto-Bäcklund transformation between two SESCOs with different degrees of sources, have been constructed and can be used to obtain N-soliton solution [11-13]. But in (2+1)-dimensional case, less results to the SESCOs have been obtained. The KP equation with self-consistent sources (KPESCS) arose in some physical models describing the interaction of long and short waves and the soliton solution of it was first found by Mel'nikov [14, 15]. However, since the explicit time part of the Lax representation of the KPESCS was not found, the method to solve the KPESCS by inverse scattering transformation in [14, 15] was quite complicated. Recently, N-soliton solution of the KPESCS was obtained by Hirota method in [16]. In this paper, we treat the constrained KP hierarchy as the stationary equations of the KP hierarchy with self-consistent sources. This gives a natural way to find the Lax pair for the KPESCS. Using the conjugate Lax pairs, we construct the generalized binary Darboux transformation with arbitrary functions in time t for the KPESCS. In contrast with the binary Darboux transformation for the KP equation which offers Bäcklund transformation, this transformation provides a non-auto-Bäcklund transformation between two KPESCSs with different degrees of sources. Some interesting solutions of KPESCS such as soliton solution, lump solution and mixture solution of exponential and rational solutions are obtained by this generalized Darboux transformation.

The paper will be organized as follows. We recall some facts about the Darboux transformation for the KP equation in the next section. In section 3, through the pseudo-differential operator (PDO) formalism we reveal the relation between the KP hierarchy with self-consistent sources and the constrained KP hierarchy. Then the conjugate Lax pairs of the KP hierarchy with self-consistent sources can be obtained naturally. Using the conjugate Lax pairs, we can construct the generalized Darboux transformations with arbitrary functions in time for KPESCS and find some interesting solutions of the KPESCS. In Section 4, the N-times repeated generalized Darboux transformation will be constructed by which the N-soliton solution and some other solutions can be obtained.

2 The Darboux transformation for the KP equation

We first give a simple description of the KP hierarchy within the framework of the Sato theory (see [20, 21]). Let us consider the following pseudo-differential operator(PDO)

$$L = \partial + u_0\partial^{-1} + u_1\partial^{-2} + \dots, \quad (2.1)$$

where ∂ denotes $\frac{\partial}{\partial x}$, and $u_j, j = 0, 1, \dots$ are functions. Denote $B_m = (L^m)_+$ for $\forall m \in \mathbb{N}$ where $(L^m)_+$ represents the projection of L^m to its pure differential part. Then the KP hierarchy has the following Lax representation (zero-curvature representation)

$$(B_n)_{t_k} - (B_k)_{t_n} + [B_n, B_k] = 0, \quad n, k \geq 2. \quad (2.2)$$

The equation (2.2) has a pair of conjugate Lax pairs as follows

$$\psi_{t_k}^- = B_k \psi^-, \quad (2.3a)$$

$$\psi_{t_n}^- = B_n \psi^-, \quad (2.3b)$$

and

$$\psi_{t_k}^+ = -(B_k)^* \psi^+, \quad (2.4a)$$

$$\psi_{t_n}^+ = -(B_n)^* \psi^+. \quad (2.4b)$$

When $k = 2, n = 3$ and under the following transformation

$$u = 2u_0, \quad t = -\frac{1}{4}t_3, \quad y = \alpha t_2, \quad (2.5)$$

we get the simplest and most important equation in the hierarchy (2.2), the KP equation

$$(u_t + 6uu_x + u_{xxx})_x + 3\alpha^2 u_{yy} = 0, \quad (2.6)$$

which appears in physical application in two different forms with $\alpha = 1$ and $\alpha = i$, usually referred to as the KPI and KPII equation [17]. Under the transformation (2.5), we obtain the conjugate Lax pairs of (2.6) respectively from (2.4) and (2.3) as follows

$$\alpha \psi_y^+ = -\psi_{xx}^+ - u \psi^+, \quad (2.7a)$$

$$\psi_t^+ = A^+(u) \psi^+, \quad A^+(u) = -4\partial^3 - 6u\partial - 3(u_x - \alpha \partial^{-1} u_y), \quad (2.7b)$$

and

$$\alpha \psi_y^- = \psi_{xx}^- + u \psi^-, \quad (2.8a)$$

$$\psi_t^- = A^-(u) \psi^-, \quad A^-(u) = -4\partial^3 - 6u\partial - 3(u_x + \alpha \partial^{-1} u_y). \quad (2.8b)$$

From (2.7) and (2.8), we can construct three types of Darboux transformations for the KP equation (2.6).

(1) The forward Darboux transformation for the KP equation.

Assume that u be a solution of the KP equation (2.6) and denote a fixed solution of (2.7) by $\psi_1^+ = \psi_1^+(x, y, t)$. The forward Darboux transformation for (2.7) is given by [19]

$$\psi^+[+1] = \psi_x^+ - \frac{\psi_{1,x}^+}{\psi_1^+} \psi^+, \quad (2.9a)$$

$$u[+1] = u + 2\partial^2 \ln \psi_1^+. \quad (2.9b)$$

So $u[+1]$ is a new solution of the KP equation (2.6). Substituting (2.9) into (2.7b), we have

$$\begin{aligned} A^+(u[+1])\psi^+[+1] &= [\psi_x^+ - \frac{\psi_{1,x}^+}{\psi_1^+} \psi^+]_t \\ &= \psi_{xt}^+ - (\frac{\psi_{1,x}^+}{\psi_1^+} \psi^+)_t \\ &= (A^+(u)\psi^+)_x - (\frac{A^+(u)\psi_1^+}{\psi_1^+})_x \psi^+ - (\frac{\psi_{1,x}^+}{\psi_1^+}) A^+(u)\psi^+. \end{aligned} \quad (2.10)$$

(2) The backward Darboux transformation for the KP equation.

Assume that u be a solution of the KP equation (2.6) and denote a fixed solution of (2.8) by $\psi_2^- = \psi_2^-(x, y, t)$. The backward Darboux transformation for the system (2.7) is defined by

$$\psi^+[-1] = \frac{a_1 + \int \psi^+ \psi_2^- dx}{\psi_2^-}, \quad (2.11a)$$

$$u[-1] = u + 2\partial^2 \ln \psi_2^-, \quad (2.11b)$$

where a_1 is an arbitrary constant. We point out that throughout the paper, the integral operation $\int f_1 f_2 dx$ such as $\int \psi^+ \psi_2^- dx$ here means $\int_{-\infty}^x f_1 f_2 dx$ or $-\int_x^{\infty} f_1 f_2 dx$ and contains no arbitrary function of y and t , only numerical constant if we impose some suitable boundary condition on the integrand functions f_1 and f_2 at $x = -\infty$ or $x = \infty$. For arbitrariness of the constants in the Darboux transformations such as a_1 here, in our computation later, the integral constants are taken to be zero.

Substituting (2.11) into (2.7b), we get the following equality

$$\begin{aligned} A^+(u[-1])\psi^+[-1] &= (\frac{a_1 + \int \psi^+ \psi_2^- dx}{\psi_2^-})_t \\ &= -\frac{A^-(u)\psi_2^-}{(\psi_2^-)^2} (a_1 + \int \psi^+ \psi_2^- dx) + \frac{\int [\psi^+ (A^-(u)\psi_2^-) + \psi_2^- (A^+(u)\psi^+)] dx}{\psi_2^-}. \end{aligned} \quad (2.12)$$

(3) The binary Darboux transformation for the KP equation.

When the backward Darboux transformation and forward Darboux transformation are applied consecutively to the system (2.7), we can get the binary DT as follows

$$\psi^+[-1, +1] = \psi^+[-1]_x - \frac{\psi_1^+[-1]_x}{\psi_1^+[-1]} \psi^+[-1] = \psi^+ - \frac{\psi_1^+(a_1 + \int \psi^+ \psi_2^- dx)}{a_2 + \int \psi_2^- \psi_1^+ dx}, \quad (2.13a)$$

$$u[-1, +1] = u[-1] + 2\partial^2 \ln \psi_1^+[-1] = u + 2\partial^2 \ln(a_2 + \int \psi_2^- \psi_1^+ dx), \quad (2.13b)$$

where $\psi_1^+[-1] = \frac{a_2 + \int \psi_1^+ \psi_2^- dx}{\psi_2^-}$, a_j , $j = 1, 2$ are arbitrary constants.

Remark: To system (2.8), we can define the forward DT, backward DT and binary DT for it analogously. We list the results here:

Forward DT of (2.8),

$$\psi^- \rightarrow \psi^-[-1] = \psi_x^- - \frac{\psi_{2,x}^-}{\psi_2^-} \psi^-, \quad (2.14a)$$

$$u \rightarrow u[-1] = u + 2\partial^2 \ln \psi_2^-, \quad (2.14b)$$

backward DT of (2.8),

$$\psi^- \rightarrow \psi^-[+1] = \frac{a_3 + \int \psi^- \psi_1^+ dx}{\psi_1^+}, \quad (2.15a)$$

$$u \rightarrow u[+1] = u + 2\partial^2 \ln \psi_1^+, \quad (2.15b)$$

Binary DT of (2.8),

$$\psi^- \rightarrow \psi^-[+1, -1] = \psi^- - \frac{\psi_1^+(a_3 + \int \psi^- \psi_1^+ dx)}{a_2 + \int \psi_2^- \psi_1^+ dx}, \quad (2.16a)$$

$$u \rightarrow u[+1, -1] = u + 2\partial^2 \ln(a_2 + \int \psi_1^+ \psi_2^- dx), \quad (2.16b)$$

where $\psi_2^-[+1] = \frac{a_2 + \int \psi_1^+ \psi_2^- dx}{\psi_1^+}$, a_j , $j = 2, 3$ are arbitrary constants.

3 The Darboux transformation for the KP equation with self-consistent sources

We first recall the constrained KP hierarchy (more details can be referred to [22-26] and the references therein). For the pseudo-differential operator L given by (2.1), if we impose the condition that $(q_j)_{t_k} = B_k q_j$, $(r_j)_{t_k} = -B_k^* q_j$, $j = 1, \dots, m$, $k \geq 2$, and make the constraint

$$(L^n)_- = \sum_{j=1}^m q_j \partial^{-1} r_j, \quad \text{or} \quad L^n = B_n + \sum_{j=1}^m q_j \partial^{-1} r_j,$$

the n-constrained KP hierarchy is defined as follows

$$(L^n)_{t_k} = [B_k, L^n], \quad (3.1a)$$

$$(q_j)_{t_k} = B_k q_j, \quad (3.1b)$$

$$(r_j)_{t_k} = -B_k^* r_j, \quad j = 1, \dots, m. \quad (3.1c)$$

If the term $(B_k)_{t_n}$ is added to the right hand side of equation (3.1a), we get the KP hierarchy with self-consistent sources as

$$(B_k)_{t_n} - (L^n)_{t_k} + [B_k, L^n] = 0, \quad (3.2a)$$

$$(q_j)_{t_k} = B_k q_j, \quad (3.2b)$$

$$(r_j)_{t_k} = -B_k^* r_j, \quad j = 1, \dots, m. \quad (3.2c)$$

If the variable " t_n " is viewed as the evolution variable, the n-constrained KP hierarchy (3.1) may be considered as the stationary case (i.e. $(B_k)_{t_n} = 0$) of the KP hierarchy with self-consistent sources (3.2). Under the condition (3.2b) and (3.2c), we naturally get the conjugate Lax pairs for (3.2a) as

$$\psi_{t_k}^- = B_k \psi^-, \quad (3.3a)$$

$$\psi_{t_n}^- = L^n \psi^- = B_n \psi^- + \sum_{j=1}^m q_j \int r_j \psi^- dx, \quad (3.3b)$$

and

$$\psi_{t_k}^+ = -B_k^* \psi^+, \quad (3.4a)$$

$$\psi_{t_n}^+ = -(L^n)^* \psi^+ = -B_n^* \psi^+ + \sum_{j=1}^m r_j \int q_j \psi^+ dx. \quad (3.4b)$$

When $k = 2, n = 3$, under the transformation (2.5) and $\Psi_j = -r_j, \Phi_j = q_j$, we will get the KPESCS [14] from (3.2) as follows

$$[u_t + 6uu_x + u_{xxx} + 8(\sum_{j=1}^m \Phi_j \Psi_j)_x]_x + 3\alpha^2 u_{yy} = 0, \quad (3.5a)$$

$$\alpha \Phi_{j,y} = \Phi_{j,xx} + u \Phi_j, \quad (3.5b)$$

$$\alpha \Psi_{j,y} = -\Psi_{j,xx} - u \Psi_j, \quad j = 1, \dots, m, \quad (3.5c)$$

and its conjugate Lax pairs from (3.4) and (3.3) respectively as follows

$$\alpha \psi_y^+ = -\psi_{xx}^+ - u \psi^+, \quad (3.6a)$$

$$\psi_t^+ = A^+(u) \psi^+ + T_m^+(\Psi, \Phi) \psi^+, \quad T_m^+(\Psi, \Phi) \psi^+ = 4 \sum_{j=1}^m \Psi_j \int \Phi_j \psi^+ dx, \quad (3.6b)$$

and

$$\alpha\psi_y^- = \psi_{xx}^- + u\psi^-, \quad (3.7a)$$

$$\psi_t^- = A^-(u)\psi^- + T_m^-(\Psi, \Phi)\psi^-, \quad T_m^-(\Psi, \Phi)\psi^- = 4 \sum_{j=1}^m \Phi_j \int \Psi_j \psi^- dx. \quad (3.7b)$$

From the conjugate Lax pairs for KPESCS (3.5), we can also construct the forward DT, backward DT and binary DT for it.

Theorem 3.1 *Assume $u, \Phi_1, \dots, \Phi_m, \Psi_1, \dots, \Psi_m$ be the solution of the KPESCS (3.5) and ψ_1^+ satisfies (3.6), then the forward Darboux transformation for (3.6) can be defined by*

$$\psi^+[+1] = \psi_x^+ - \frac{\psi_{1,x}^+}{\psi_1^+} \psi^+, \quad (3.8a)$$

$$u[+1] = u + 2\partial^2 \ln \psi_1^+, \quad (3.8b)$$

$$\Phi_j[+1] = -\frac{\int \psi_1^+ \Phi_j dx}{\psi_1^+}, \quad (3.8c)$$

$$\Psi_j[+1] = \Psi_{j,x} - \frac{\psi_{1,x}^+}{\psi_1^+} \Psi_j, \quad j = 1, \dots, m. \quad (3.8d)$$

Namely, $u[+1], \psi^+[+1], \Psi_j[+1], \Phi_j[+1], j = 1, \dots, m$, satisfy (3.5) and (3.6).

Proof: Based on the results in the previous section, it is obvious that (3.5b), (3.5c) and (3.6a) hold under the transformation (3.8). So we only need to prove (3.6b), i.e., the following equality,

$$\begin{aligned} \psi^+[+1]_t &= [\psi_x^+ - \frac{\psi_{1,x}^+}{\psi_1^+} \psi^+]_t \\ &= (A^+(u)\psi^+ + T_m^+(\Psi, \Phi)\psi^+)_x - (\frac{A^+(u)\psi_1^+ + T_m^+(\Psi, \Phi)\psi_1^+}{\psi_1^+})_x \psi^+ \\ &\quad - \frac{\psi_{1,x}^+}{\psi_1^+} (A^+(u)\psi^+ + T_m^+(\Psi, \Phi)\psi^+) \\ &= (A^+(u)\psi^+)_x - (\frac{A^+(u)\psi_1^+}{\psi_1^+})_x \psi^+ - \frac{\psi_{1,x}^+}{\psi_1^+} A^+(u)\psi^+ \\ &\quad + (T_m^+(\Psi, \Phi)\psi^+)_x - (\frac{T_m^+(\Psi, \Phi)\psi_1^+}{\psi_1^+})_x \psi^+ - \frac{\psi_{1,x}^+}{\psi_1^+} T_m^+(\Psi, \Phi)\psi^+ \\ &= A^+(u[+1])\psi^+[+1] + T_m^+(\Psi[+1], \Phi[+1])\psi^+[+1]. \end{aligned} \quad (3.9)$$

By simple computation, we can prove that equality (2.10) still holds now, so we only need to check the terms containing Φ_j, Ψ_j in the equality (3.9), i.e., to check the following equality

$$(T_m^+(\Psi, \Phi)\psi^+)_x - \left(\frac{T_m^+(\Psi, \Phi)\psi_1^+}{\psi_1^+}\right)_x \psi^+ - \frac{\psi_{1,x}^+}{\psi_1^+} T_m^+(\Psi, \Phi)\psi^+ = T_m^+(\Psi[+1], \Phi[+1])\psi^+[+1]. \quad (3.10)$$

In fact, we have

$$\begin{aligned} & \text{the l.h.s. of (3.10)} \\ &= 4[\sum_{j=1}^m \Psi_j \int \Phi_j \psi^+ dx]_x - 4[\frac{\sum_{j=1}^m \Psi_j \int \Phi_j \psi_1^+ dx}{\psi_1^+}]_x \psi^+ - 4\frac{\psi_{1,x}^+}{\psi_1^+} \sum_{j=1}^m \Psi_j \int \Phi_j \psi^+ dx \\ &= 4 \sum_{j=1}^m [\Psi_{j,x} \int \Phi_j \psi^+ dx + \Psi_j \Phi_j \psi^+] - 4\frac{\psi_{1,x}^+}{\psi_1^+} \sum_{j=1}^m \Psi_j \int \Phi_j \psi^+ dx \\ &\quad - 4\frac{\sum_{j=1}^m [(\Psi_{j,x} \int \Phi_j \psi_1^+ dx + \Psi_j \Phi_j \psi_1^+) \psi_1^+ - (\Psi_j \int \Phi_j \psi_1^+ dx) \psi_{1,x}^+]}{(\psi_1^+)^2} \psi^+ \\ &= 4 \sum_{j=1}^m [\Psi_{j,x} \int \Phi_j \psi^+ dx - \frac{\psi^+ \Psi_{j,x}}{\psi_1^+} \int \Phi_j \psi_1^+ dx + \Psi_j \psi^+ \psi_{1,x}^+ \frac{\int \Phi_j \psi_1^+ dx}{(\psi_1^+)^2} - \frac{\psi_{1,x}^+}{\psi_1^+} \Psi_j \int \Phi_j \psi^+ dx] \\ &= 4 \sum_{j=1}^m [(\Psi_{j,x} - \frac{\psi_{1,x}^+}{\psi_1^+} \Psi_j) \int \Phi_j \psi^+ dx - \frac{\psi^+ \int \Phi_j \psi_1^+ dx}{\psi_1^+} (\Psi_{j,x} - \frac{\psi_{1,x}^+}{\psi_1^+} \Psi_j)] \\ &= 4 \sum_{j=1}^m \Psi_j[+1] (\int \Phi_j \psi^+ dx - \frac{\psi^+}{\psi_1^+} \int \Phi_j \psi_1^+ dx) \end{aligned} \quad (3.11)$$

$$\begin{aligned} \text{the r.h.s. of (3.10)} &= 4 \sum_{j=1}^m \Psi_j[+1] \int \Phi_j[+1] \psi^+[+1] dx \\ &= 4 \sum_{j=1}^m \Psi_j[+1] \int [-\frac{\int \psi_1^+ \Phi_j dx}{\psi_1^+} (\psi_x^+ - \frac{\psi_{1,x}^+}{\psi_1^+} \psi^+)] dx \\ &= -4 \sum_{j=1}^m \Psi_j[+1] \int [(\int \psi_1^+ \Phi_j dx) d(\frac{\psi^+}{\psi_1^+})] dx \\ &= 4 \sum_{j=1}^m \Psi_j[+1] [\int \Phi_j \psi^+ dx - \frac{\psi^+}{\psi_1^+} \int \Phi_j \psi_1^+ dx] \\ &= \text{the l.h.s. of (3.10)} \end{aligned} \quad (3.12)$$

This completes the proof.

Theorem 3.2 Assume $u, \Phi_1, \dots, \Phi_m, \Psi_1, \dots, \Psi_m$ be a solution of the KPESCS (3.5) and ψ_2^- satisfies (3.7), then the backward DT for (3.6) is defined by

$$\psi^+[-1] = \frac{a_4 + \int \psi_2^- \psi^+ dx}{\psi_2^-}, \quad (3.13a)$$

$$u[-1] = u + 2\partial^2 \ln \psi_2^-, \quad (3.13b)$$

$$\Phi_j[-1] = \Phi_{j,x} - \frac{\psi_{2,x}^-}{\psi_2^-} \Phi_j, \quad (3.13c)$$

$$\Psi_j[-1] = -\frac{\int \psi_2^- \Psi_j dx}{\psi_2^-}, \quad j = 1, \dots, m, \quad (3.13d)$$

where a_4 is an arbitrary constant, namely $u[-1]$, $\psi^+[-1]$, $\Psi_j[-1]$, $\Phi_j[-1]$, $j = 1, \dots, m$, satisfy (3.5) and (3.6).

Proof: It is obvious that (3.5b), (3.5c) and (3.6a) hold under the transformation (3.13). So we only need to prove (3.6b), i.e., to prove the following equality

$$\begin{aligned} & \psi^+[-1]_t \\ = & \left[\frac{a_4 + \int \psi_2^- \psi^+ dx}{\psi_2^-} \right]_t \\ = & -\frac{A^-(u)\psi_2^-}{(\psi_2^-)^2} (a_4 + \int \psi_2^- \psi^+ dx) + \frac{\int [(A^-(u)\psi_2^-)\psi^+ + (A^+(u)\psi^+)\psi_2^-] dx}{\psi_2^-} \\ & -\frac{T_m^-(\Phi, \Psi)\psi_2^-}{(\psi_2^-)^2} (a_4 + \int \psi_2^- \psi^+ dx) + \frac{\int [(T_m^-(\Phi, \Psi)\psi_2^-)\psi^+ + (T_m^+(\Psi, \Phi)\psi^+)\psi_2^-] dx}{\psi_2^-} \\ = & A^+(u[-1])\psi^+[-1] + T_m^+(\Psi[-1], \Phi[-1])\psi^+[-1]. \end{aligned} \quad (3.14)$$

Similarly, using equality (2.12), we only need to check the terms containing Ψ_j, Φ_j , i.e.,

$$\begin{aligned} & -\frac{T_m^-(\Psi, \Phi)\psi_2^-}{(\psi_2^-)^2} (a_4 + \int \psi_2^- \psi^+ dx) + \frac{\int [(T_m^-(\Psi, \Phi)\psi_2^-)\psi^+ + (T_m^+(\Psi, \Phi)\psi^+)\psi_2^-] dx}{\psi_2^-} \\ = & T_m^+(\Psi[-1], \Phi[-1])\psi^+[-1], \end{aligned} \quad (3.15)$$

i.e.,

$$\begin{aligned} & 4 \sum_{j=1}^m \left\{ -\frac{\Phi_j \int \psi_2^- \Psi_j dx}{(\psi_2^-)^2} (a_4 + \int \psi_2^- \psi^+ dx) + \frac{\int [\psi^+ \Phi_j \int (\psi_2^- \Psi_j) + \psi_2^- \Psi_j \int (\psi^+ \Phi_j)] dx}{\psi_2^-} \right\} \\ = & 4 \sum_{j=1}^m \Psi_j[-1] \int \Phi_j[-1] \psi^+[-1] dx. \end{aligned} \quad (3.16)$$

In fact, we have

$$\begin{aligned} \text{the l.h.s. of (3.16)} &= 4 \sum_{j=1}^m \left[-\frac{\Phi_j \int \psi_2^- \Psi_j dx}{(\psi_2^-)^2} (a_4 + \int \psi_2^- \psi^+ dx) + \frac{\int \psi^+ \Phi_j dx \int \psi_2^- \Psi_j dx}{\psi_2^-} \right] \\ &= 4 \sum_{j=1}^m \frac{\int \Psi_j \psi_2^- dx}{\psi_2^-} \left[-\frac{\Phi_j}{\psi_2^-} (a_4 + \int \psi_2^- \psi^+ dx) + \int \psi^+ \Phi_j dx \right] \\ &= -4 \sum_{j=1}^m \frac{\int \Psi_j \psi_2^- dx}{\psi_2^-} \left\{ \int [(a_4 + \int \psi_2^- \psi^+ dx) d(\frac{\Phi_j}{\psi_2^-})] dx \right\} \\ &= -4 \sum_{j=1}^m \frac{\int \Psi_j \psi_2^- dx}{\psi_2^-} \int \left[\frac{(a_4 + \int \psi_2^- \psi^+ dx)}{\psi_2^-} \frac{(\Phi_{j,x} \psi_2^- - \Phi_j \psi_{2,x}^-)}{\psi_2^-} \right] dx \\ &= 4 \sum_{j=1}^m \Psi_j[-1] \int \psi^+[-1] \Phi_j[-1] dx \\ &= \text{the r.h.s. of (3.16)}. \end{aligned} \quad (3.17)$$

This completes the proof.

From Theorem 3.1 and Theorem 3.2, we can obtain the binary Darboux transformation for the system (3.6) by choosing $a_4 = 0$, $\psi_1^+[-1] = \frac{C + \int \psi_1^+ \psi_2^- dx}{\psi_2^-}$ (C is a constant) as follows

$$\psi^+[-1, +1] = \psi^+ - \frac{\psi_1^+ \int \psi_2^- \psi^+ dx}{C + \int \psi_1^+ \psi_2^- dx}, \quad (3.18a)$$

$$u[-1, +1] = u + 2\partial^2 \ln(C + \int \psi_2^- \psi_1^+ dx), \quad (3.18b)$$

$$\Phi_j[-1, +1] = \Phi_j - \frac{\psi_2^- \int \psi_1^+ \Phi_j dx}{C + \int \psi_2^- \psi_1^+ dx}, \quad (3.18c)$$

$$\Psi_j[-1, +1] = \Psi_j - \frac{\psi_1^+ \int \psi_2^- \Psi_j dx}{C + \int \psi_2^- \psi_1^+ dx}, \quad j = 1, \dots, m. \quad (3.18d)$$

Substituting (3.18) into equation (3.6b) gives the equality

$$\psi^+[-1, +1]_t = A^+(u[-1, +1])\psi^+[-1, +1] + T_m^+(\Psi[-1, +1], \Phi[-1, +1])\psi^+[-1, +1]. \quad (3.19)$$

Both sides of the equality (3.19) are polynomials w.r.t. the term $(C + \int \psi_2^- \psi_1^+ dx)^{-1}$. For example, the left hand side of (3.19) is a polynomial of order 2 w.r.t. $(C + \int \psi_2^- \psi_1^+ dx)^{-1}$ as follows

$$\begin{aligned} \psi^+[-1, +1]_t &= [\psi^+ - \frac{\psi_1^+ \int \psi_2^- \psi^+ dx}{C + \int \psi_1^+ \psi_2^- dx}]_t \\ &= \psi_t^+ - \frac{\psi_{1,t}^+ (\int \psi_2^- \psi^+ dx)}{C + \int \psi_1^+ \psi_2^- dx} - \frac{\psi_1^+ (\int \psi_2^- \psi^+ dx)_t}{C + \int \psi_1^+ \psi_2^- dx} + \frac{\psi_1^+ (\int \psi_2^- \psi^+ dx) (\int \psi_2^- \psi^+ dx)_t}{(C + \int \psi_1^+ \psi_2^- dx)^2}. \\ &\equiv \sum_{j=0}^2 L_j [C + \int \psi_2^- \psi_1^+ dx]^{-j}. \end{aligned} \quad (3.20)$$

By a tedious computation, the right hand side of (3.19) is displayed to be a polynomial of the order less than 4 w.r.t. $(C + \int \psi_2^- \psi_1^+ dx)^{-1}$. We denote it by $\sum_{j=0}^4 R_j [C + \int \psi_2^- \psi_1^+ dx]^{-j}$. Since (3.19) holds for arbitrary constant C and L_j, R_j never contain C , we have

$$L_j = R_j, \quad j = 0, 1, 2 \quad \text{and} \quad R_j = 0, \quad j = 3, 4.$$

If we replace C in (3.18) by $C(t)$, an arbitrary function in t , and substitute equations (3.18) into both sides of the equations of (3.6) again, we will find that (3.6a) is also covariant w.r.t. (3.18). The right hand side of (3.19) turns to be $\sum_{j=0}^2 R_j [C(t) + \int \psi_2^- \psi_1^+ dx]^{-j}$, but the left hand side of (3.19) does not equal to $\sum_{j=0}^2 L_j [C(t) + \int \psi_2^- \psi_1^+ dx]^{-j}$ any more. In other words, (3.6b) does not covariant w.r.t. (3.18) any longer when C replaced by $C(t)$. In fact we have the following theorem.

Theorem 3.3 Given $u, \Psi_1, \dots, \Psi_m, \Phi_1, \dots, \Phi_m$ a solution of the KPESCS (3.5) and let ψ_1^+ and ψ_2^- be solutions of the system (3.6) and (3.7) respectively, then the transformation with $C(t)$, an arbitrary function in t defined by

$$\psi^+[-1, +1] = \psi^+ - \frac{\psi_1^+ \int \psi_2^- \psi^+ dx}{C(t) + \int \psi_1^+ \psi_2^- dx}, \quad (3.21a)$$

$$u[-1, +1] = u + 2\partial^2 \ln(C(t) + \int \psi_2^- \psi_1^+ dx), \quad (3.21b)$$

$$\Phi_j[-1, +1] = \Phi_j - \frac{\psi_2^- \int \psi_1^+ \Phi_j dx}{C(t) + \int \psi_2^- \psi_1^+ dx}, \quad (3.21c)$$

$$\Psi_j[-1, +1] = \Psi_j - \frac{\psi_1^+ \int \psi_2^- \Psi_j dx}{C(t) + \int \psi_2^- \psi_1^+ dx}, \quad j = 1, \dots, m, \quad (3.21d)$$

and

$$\Psi_{m+1}[-1, +1] = \frac{1}{2} \frac{\sqrt{\dot{C}(t)} \psi_1^+}{C(t) + \int \psi_2^- \psi_1^+ dx}, \quad \Phi_{m+1}[-1, +1] = \frac{1}{2} \frac{\sqrt{\dot{C}(t)} \psi_2^-}{C(t) + \int \psi_2^- \psi_1^+ dx}, \quad (3.21e)$$

transforms (3.5b), (3.5c) and (3.6) respectively into

$$\alpha \Phi_j[-1, +1]_y = \Phi_j[-1, +1]_{xx} + u[-1, +1] \Phi_j[-1, +1], \quad (3.22a)$$

$$\alpha \Psi_j[-1, +1]_y = -\Psi_j[-1, +1]_{xx} - u[-1, +1] \Psi_j[-1, +1], \quad j = 1, \dots, m+1, \quad (3.22b)$$

$$\alpha \psi^+[-1, +1]_y = -\psi^+[-1, +1]_{xx} - u[-1, +1] \psi^+[-1, +1], \quad (3.22c)$$

$$\psi^+[-1, +1]_t = A^+(u[-1, +1]) \psi^+[-1, +1] + T_{m+1}^+(\Psi[-1, +1], \Phi[-1, +1]) \psi^+[-1, +1]. \quad (3.22d)$$

So $u[-1, +1], \Psi_1[-1, +1], \dots, \Psi_{m+1}[-1, +1], \Phi_1[-1, +1], \dots, \Phi_{m+1}[-1, +1]$ is a solution of the KPESCS (3.5) with degree $m+1$.

Proof: Equations (3.22a), (3.22b) and (3.22c) hold obviously. We only need to prove (3.22d). Substituting (3.21a) into the left hand side of (3.22d), we have

$$\begin{aligned} \psi^+[-1, +1]_t &= [\psi^+ - \frac{\psi_1^+ \int \psi_2^- \psi^+ dx}{C(t) + \int \psi_1^+ \psi_2^- dx}]_t \\ &= \sum_{j=0}^2 L_j [C(t) + \int \psi_2^- \psi_1^+ dx]^{-j} + \frac{\dot{C}(t) \psi_1^+ \int \psi_2^- \psi^+ dx}{(C(t) + \int \psi_1^+ \psi_2^- dx)^2} \\ &= \sum_{j=0}^2 R_j [C(t) + \int \psi_2^- \psi_1^+ dx]^{-j} + \frac{\dot{C}(t) \psi_1^+ \int \psi_2^- \psi^+ dx}{(C(t) + \int \psi_1^+ \psi_2^- dx)^2} \\ &= A^+(u[-1, +1]) \psi^+[-1, +1] + T_m^+(\Psi[-1, +1], \Phi[-1, +1]) \psi^+[-1, +1] \\ &\quad + \frac{\dot{C}(t) \psi_1^+ \int \psi_2^- \psi^+ dx}{(C(t) + \int \psi_1^+ \psi_2^- dx)^2}. \end{aligned} \quad (3.23)$$

So we only need to prove

$$4\Psi_{m+1}[-1, +1] \int \Phi_{m+1}[-1, +1] \psi^+[-1, +1] dx = \frac{\dot{C}(t) \psi_1^+ \int \psi^+ \psi_2^- dx}{(C(t) + \int \psi_1^+ \psi_2^- dx)^2}. \quad (3.24)$$

In fact, we have

$$\begin{aligned} \text{the l.h.s of (3.24)} &= 4 \times \frac{1}{2} \frac{\sqrt{\dot{C}(t)} \psi_1^+}{C(t) + \int \psi_1^+ \psi_2^- dx} \int \left[\left(\psi^+ - \frac{\psi_1^+ \int \psi^+ \psi_2^- dx}{C(t) + \int \psi_1^+ \psi_2^- dx} \right) \times \frac{1}{2} \frac{\sqrt{\dot{C}(t)} \psi_2^-}{C(t) + \int \psi_2^- \psi_1^+ dx} \right] dx \\ &= \frac{\dot{C}(t) \psi_1^+}{C(t) + \int \psi_1^+ \psi_2^- dx} \int \left[\left(\psi^+ - \frac{\psi_1^+ \int \psi^+ \psi_2^- dx}{C(t) + \int \psi_1^+ \psi_2^- dx} \right) \frac{\psi_2^-}{C(t) + \int \psi_1^+ \psi_2^- dx} \right] dx \\ &= \frac{\dot{C}(t) \psi_1^+}{C(t) + \int \psi_1^+ \psi_2^- dx} \int \frac{\psi^+ \psi_2^- (C(t) + \int \psi_1^+ \psi_2^- dx) - \psi_1^+ \psi_2^- (\int \psi^+ \psi_2^- dx)}{(C(t) + \int \psi_1^+ \psi_2^- dx)^2} dx, \\ &= \frac{\dot{C}(t) \psi_1^+}{C(t) + \int \psi_1^+ \psi_2^- dx} \int d \left(\frac{\int \psi^+ \psi_2^- dx}{C(t) + \int \psi_1^+ \psi_2^- dx} \right) \\ &= \frac{\dot{C}(t) \psi_1^+}{C(t) + \int \psi_1^+ \psi_2^- dx} \frac{\int \psi^+ \psi_2^- dx}{C(t) + \int \psi_1^+ \psi_2^- dx} \\ &= \text{the r.h.s of (3.24)} \end{aligned} \quad (3.25)$$

This completes the proof.

Remark:

(1). For system (3.7), the DTs described in this section can also be constructed, we omit the results here.

(2). (3.21b)-(3.21e) offer us a non-auto-Bäcklund transformation between two KPESCSs with degrees m and $m + 1$ respectively. This DT can be used to construct solutions of the KPESCS (3.5).

Example 1: single-soliton solution of KPI equation with self-consistent sources (KPIESCS).

If we set $\alpha = i$ in equation (3.5), we get the KPIESCS [15]

$$[u_t + 6uu_x + u_{xxx} + 8 \sum_{j=1}^m (\Phi_j \Psi_j)_x]_x - 3u_{yy} = 0, \quad (3.26a)$$

$$i\Phi_{j,y} = \Phi_{j,xx} + u\Phi_j, \quad (3.26b)$$

$$i\Psi_{j,y} = -\Psi_{j,xx} - u\Psi_j, \quad j = 1, \dots, m. \quad (3.26c)$$

We take $u = 0$ as the initial solution of (3.26) with $m = 0$ and let

$$\psi_1^+ = e^{k_1 x + i k_1^2 y - 4 k_1^3 t}, \quad \psi_2^- = e^{k_2 x - i k_2^2 y - 4 k_2^3 t}, \quad k_1 = \mu + i\nu, \quad k_2 = \mu - i\nu, \quad \mu, \nu \in \mathbb{R}, \quad (3.27a)$$

$$C(t) = \frac{1}{k_1 + k_2} e^{2\beta(t)} = \frac{1}{2\mu} e^{2\beta(t)}, \quad (3.27b)$$

where $\beta(t)$ is an arbitrary function in t . Then $\psi_1^+ = e^{\omega + \theta}$, $\psi_2^- = e^{-\omega + \theta}$ where

$$\omega = -i\nu x + i(\mu^2 - \nu^2)y + 4i(3\mu^2\nu - \nu^3)t, \quad \theta = \mu x + 2\mu\nu y - 4(\mu^3 - 3\mu\nu^2)t.$$

The single-soliton solution of the KPIESCS (3.26) with $m = 1$ is given by

$$u[-1, +1] = 2\partial^2 \ln(C(t) + \int \psi_1^+ \psi_2^- dx) = 2\partial^2 \ln\left(\frac{1}{2\mu} e^{2\beta(t)} + \frac{1}{2\mu} e^{2\theta(t)}\right) = 2\mu^2 \operatorname{sech}^2(\theta - \beta(t)), \quad (3.28a)$$

$$\Phi_1[-1, +1] = \frac{1}{2} \frac{\sqrt{\dot{C}(t)} \psi_2^-}{C(t) + \frac{1}{2\mu} e^{2\theta(t)}} = a e^{-\omega} \operatorname{sech}(\theta - \beta(t)), \quad a = \frac{1}{2} \sqrt{\mu \dot{\beta}(t)} \quad (3.28b)$$

$$\Psi_1[-1, +1] = \overline{\Phi_1[-1, +1]} = a e^{\omega} \operatorname{sech}(\theta - \beta(t)). \quad (3.28c)$$

We can compare the results above with those in [15].

Case 1: If we set

$$\theta - \beta(t) = \mu x + 2\mu\nu y - 4(\mu^3 - 3\mu\nu^2)t - \beta(t) = \mu(x + 2\nu y - \tau t),$$

then

$$\beta(t) = [\tau - 4(\mu^2 - 3\nu^2)]\mu t,$$

$$\dot{\beta}(t) = \mu[\tau - 4(\mu^2 - 3\nu^2)],$$

then we have the following identity

$$a^2 = \frac{1}{4} \mu \dot{\beta}(t) = \frac{\mu^2}{4} [\tau - 4(\mu^2 - 3\nu^2)], \quad (3.29)$$

which gives the relation of the parameters a , τ , μ , ν appearing in the solution (3.28). The solution (3.28) together with (3.29) is also obtained in [15].

Case 2: More generally, if we set

$$\theta - \beta(t) = \mu(x + 2\nu y - f(t)),$$

then

$$f(t) = 4(\mu^2 - 3\nu^2)t + \frac{\beta(t)}{\mu},$$

then the relation identity of the parameters is

$$\frac{df}{dt}(t) = 4(\mu^2 - 3\nu^2) + \frac{\dot{\beta}(t)}{\mu} = 4(\mu^2 - 3\nu^2) + 4\mu^{-2}a^2, \quad (3.30)$$

which is also given in [15].

From **Case 1** and **Case 2**, we see that the solutions obtained through the generalized binary DT are consistent with those in [15], and the method used here is more natural and simple.

Example 2: lump solution (rational solution) of the KPIESCS.

We take $u = 0$, $\Phi_1 = \Psi_1 = 1$ as the initial solution of (3.26) with $m = 1$ and let

$$\psi_1^+ = (x + 2iky - 12k^2t - \frac{4}{k^2}t)e^{kx+ik^2y-4k^3t+\frac{4}{k}t}, \quad \psi_2^- = \overline{\psi_1^+}, \quad k \in \mathbb{C}.$$

Set $k = \mu + i\nu$, $\mu, \nu \in \mathbb{R}$, then

$$\psi_1^+ = e^{\theta_1+i\theta_2}(\alpha_1 + \alpha_2 i), \quad \psi_2^- = e^{\theta_1-i\theta_2}(\alpha_1 - \alpha_2 i),$$

where

$$\theta_1 = \mu x - 2\mu\nu y - 4\mu(\mu^2 - 3\nu^2)t + \frac{4\mu}{\mu^2 + \nu^2}t, \quad \theta_2 = \nu x + (\mu^2 - \nu^2)y - 4\nu(3\mu^2 - \nu^2)t - \frac{4\nu}{\mu^2 + \nu^2}t,$$

$$\alpha_1 = x - 2\nu y - 12(\mu^2 - \nu^2)t - 4\frac{\mu^2 - \nu^2}{(\mu^2 + \nu^2)^2}t, \quad \alpha_2 = 2\mu y - 24\mu\nu t + \frac{8\mu\nu}{(\mu^2 + \nu^2)^2}t.$$

By DT (3.21) with $C(t) \equiv 0$, we get the lump solution of the KPIESCS (3.26) with $m = 1$ as follows

$$u[-1, +1] = 2\partial^2 \ln \int \psi_1^+ \psi_2^- dx = 2\partial^2 \ln [\alpha_2^2 + (\alpha_1 - \frac{1}{2\mu})^2 + \frac{1}{4\mu^2}], \quad (3.31a)$$

$$\begin{aligned} \Phi_1[-1, +1] &= \Phi_1 - \frac{\psi_2^- \int \psi_1^+ \Phi_1 dx}{\int \psi_1^+ \psi_2^- dx} \\ &= 1 - 2\mu \frac{(\alpha_1^2 + \alpha_2^2)\mu(\mu^2 + \nu^2) - \alpha_1(\mu^2 - \nu^2) + 2\mu\nu\alpha_2 + [-(\alpha_1^2 + \alpha_2^2)(\nu\mu^2 + \nu^3) + (\mu^2 - \nu^2)\alpha_2 + 2\mu\nu\alpha_1]i}{(\mu^2 + \nu^2)^2[\alpha_2^2 + (\alpha_1 - \frac{1}{2\mu})^2 + \frac{1}{4\mu^2}]}, \end{aligned} \quad (3.31b)$$

$$\begin{aligned} \Psi_1[-1, +1] &= \overline{\Phi_1[-1, +1]} \\ &= 1 - 2\mu \frac{(\alpha_1^2 + \alpha_2^2)\mu(\mu^2 + \nu^2) - \alpha_1(\mu^2 - \nu^2) + 2\mu\nu\alpha_2 - [-(\alpha_1^2 + \alpha_2^2)(\nu\mu^2 + \nu^3) + (\mu^2 - \nu^2)\alpha_2 + 2\mu\nu\alpha_1]i}{(\mu^2 + \nu^2)^2[\alpha_2^2 + (\alpha_1 - \frac{1}{2\mu})^2 + \frac{1}{4\mu^2}]}. \end{aligned} \quad (3.31c)$$

Example 3: single-soliton solution of the KP II equation with self-consistent sources (KPI-IESCS).

If we set $\alpha = 1$ in equation (3.5), we get the KPIIESCS [16]

$$[u_t + 6uu_x + u_{xxx} + 8 \sum_{j=1}^m (\Phi_j \Psi_j)_x]_x + 3u_{yy} = 0, \quad (3.32a)$$

$$\Phi_{j,y} = \Phi_{j,xx} + u\Phi_j, \quad (3.32b)$$

$$\Psi_{j,y} = -\Psi_{j,xx} - u\Psi_j, \quad j = 1, \dots, m. \quad (3.32c)$$

We take $u = 0$ as the initial solution of (3.32) with $m = 0$ and let

$$\psi_1^+ = e^{k_1 x - k_1^2 y - 4k_1^3 t} = e^{\xi_1}, \quad \psi_2^- = e^{k_2 x + k_2^2 y - 4k_2^3 t} = e^{\xi_2}, \quad k_1, k_2 \in \mathbb{R}, \quad (3.33a)$$

$$C(t) = \frac{1}{k_1 + k_2} e^{2\beta(t)}, \quad (3.33b)$$

where $\beta(t)$ is an arbitrary function in t . The single-soliton solution of the KPIESCS (3.32) with $m = 1$ is given by

$$\begin{aligned} u[-1, +1] &= 2\partial^2 \ln(C(t) + \int \psi_1^+ \psi_2^- dx) = 2\partial^2 \ln(1 + e^{\xi_1 + \xi_2 - 2\beta(t)}), \\ \Phi_1[-1, +1] &= \frac{1}{2} \frac{\sqrt{\dot{C}(t)} \psi_2^-}{C(t) + \frac{1}{k_1 + k_2} e^{\xi_1 + \xi_2}} = \frac{\sqrt{2(k_1 + k_2)} \dot{\beta} e^{\xi_2 - \beta(t)}}{2(1 + e^{\xi_1 + \xi_2 - 2\beta(t)})}, \\ \Psi_1[-1, +1] &= \frac{1}{2} \frac{\sqrt{\dot{C}(t)} \psi_1^+}{C(t) + \frac{1}{k_1 + k_2} e^{\xi_1 + \xi_2}} = \frac{\sqrt{2(k_1 + k_2)} \dot{\beta} e^{\xi_1 - \beta(t)}}{2(1 + e^{\xi_1 + \xi_2 - 2\beta(t)})}. \end{aligned} \quad (3.34)$$

4 The N-times Repeated Generalized Binary Darboux transformation for the KPESCS

Assuming f_1, \dots, f_n are n arbitrary solutions of (3.6), g_1, \dots, g_n are n arbitrary solutions of (3.7), $C_1(t), \dots, C_n(t)$ are n arbitrary functions in t , we define the following Wronskians:

$$\begin{aligned} W_1(f_1, \dots, f_n; g_1, \dots, g_n; C_1(t), \dots, C_n(t)) &= \det(U_{n \times n}), \\ W_2(f_1, \dots, f_n; g_1, \dots, g_{n-1}; C_1(t), \dots, C_{n-1}(t)) &= \det(V_{n \times n}), \\ W_3(f_1, \dots, f_{n-1}; g_1, \dots, g_n; C_1(t), \dots, C_{n-1}(t)) &= \det(X_{n \times n}), \end{aligned} \quad (4.1)$$

where

$$U_{i,j} = \delta_{i,j} C_i(t) + \int g_i f_j dx, \quad i, j = 1, \dots, n, \quad (4.2a)$$

$$V_{i,j} = \delta_{i,j} C_i(t) + \int g_i f_j dx, \quad i = 1, \dots, n-1, \quad j = 1, \dots, n; \quad V_{n,j} = f_j, \quad j = 1, \dots, n, \quad (4.2b)$$

$$X_{i,j} = \delta_{i,j} C_i(t) + \int g_j f_i dx, \quad i = 1, \dots, n-1, \quad j = 1, \dots, n; \quad X_{n,j} = g_j, \quad j = 1, \dots, n. \quad (4.2c)$$

Lemma 4.1 Assume $\psi^+, \psi_1^+, \dots, \psi_N^+$ are solutions of (3.6) and $\psi^-, \psi_{N+1}^-, \dots, \psi_{N+N}^-$ are solutions of (3.7). For notation simplicity, we denote $\psi_{l+k}^+[l] = \psi_{l+k}^+[-l, +l]$, $\psi_{N+l+k}^-[l] = \psi_{N+l+k}^-[-l, +l]$, $0 \leq l \leq N$, $1 \leq k \leq N-l$. Then

$$\begin{aligned} &W_1(\psi_m^+[m-1], \dots, \psi_{m+k}^+[m-1]; \psi_{N+m}^-[m-1], \dots, \psi_{N+m+k}^-[m-1]; C_m(t), \dots, C_{m+k}(t)) \\ &= \frac{W_1(\psi_{m-1}^+[m-2], \dots, \psi_{m+k}^+[m-2]; \psi_{N+m-1}^-[m-2], \dots, \psi_{N+m+k}^-[m-2]; C_{m-1}(t), \dots, C_{m+k}(t))}{C_{m-1}(t) + \int \psi_{m-1}^+[m-2] \psi_{N+m-1}^-[m-2] dx}, \end{aligned} \quad (4.3a)$$

$$\begin{aligned}
& W_2(\psi_m^+[m-1], \dots, \psi_{m+k}^+[m-1], \psi^+[m-1]; \psi_{N+m}^-[m-1], \dots, \psi_{N+m+k}^-[m-1]; \\
& C_m(t), \dots, C_{m+k}(t)) \\
= & \frac{W_2(\psi_{m-1}^+[m-2], \dots, \psi_{m+k}^+[m-2], \psi^+[m-2]; \psi_{N+m-1}^-[m-2], \dots, \psi_{N+m+k}^-[m-2]; C_{m-1}(t), \dots, C_{m+k}(t))}{C_{m-1}(t) + \int \psi_{m-1}^+[m-2] \psi_{N+m-1}^-[m-2] dx},
\end{aligned} \tag{4.3b}$$

$$\begin{aligned}
& W_3(\psi_m^+[m-1], \dots, \psi_{m+k}^+[m-1]; \psi_{N+m}^-[m-1], \dots, \psi_{N+m+k}^-[m-1], \psi^-[m-1]; \\
& C_m(t), \dots, C_{m+k}(t)) \\
= & \frac{W_3(\psi_{m-1}^+[m-2], \dots, \psi_{m+k}^+[m-2]; \psi_{N+m-1}^-[m-2], \dots, \psi_{N+m+k}^-[m-2], \psi^-[m-2]; C_{m-1}(t), \dots, C_{m+k}(t))}{C_{m-1}(t) + \int \psi_{m-1}^+[m-2] \psi_{N+m-1}^-[m-2] dx}.
\end{aligned} \tag{4.3c}$$

Proof: According to (3.21), we have

$$\psi_{m-1+j}^+[m-1] = \psi_{m-1+j}^+[m-2] - \frac{\psi_{m-1}^+[m-2] \int \psi_{N+m-1}^-[m-2] \psi_{m-1+j}^+[m-2] dx}{C_{m-1}(t) + \int \psi_{m-1}^+[m-2] \psi_{N+m-1}^-[m-2] dx}, \tag{4.4a}$$

$$\psi_{N+m-1+i}^-[m-1] = \psi_{N+m-1+i}^-[m-2] - \frac{\psi_{N+m-1}^-[m-2] \int \psi_{m-1}^+[m-2] \psi_{N+m-1+i}^-[m-2] dx}{C_{m-1}(t) + \int \psi_{m-1}^+[m-2] \psi_{N+m-1}^-[m-2] dx}. \tag{4.4b}$$

By definition (4.2a) and equations (4.4), we have

$$\begin{aligned}
& W_1(\psi_m^+[m-1], \dots, \psi_{m+k}^+[m-1]; \psi_{N+m}^-[m-1], \dots, \psi_{N+m+k}^-[m-1]; C_m(t), \dots, C_{m+k}(t)) \\
= & \det(U_{i,j}), \quad 1 \leq i, j \leq k+1,
\end{aligned} \tag{4.5}$$

$$\begin{aligned}
& U_{i,j} \\
= & \delta_{i,j} C_{m-1+i}(t) + \int \psi_{m-1+j}^+[m-1] \psi_{N+m-1+i}^-[m-1] dx \\
= & \delta_{i,j} C_{m-1+i}(t) + \int \psi_{m-1+j}^+[m-2] \psi_{N+m-1+i}^-[m-2] dx - \frac{1}{C_{m-1}(t) + \int \psi_{m-1}^+[m-2] \psi_{N+m-1}^-[m-2] dx} \\
& (\int \psi_{m-1+j}^+[m-2] \psi_{N+m-1}^-[m-2] dx) (\int \psi_{m-1}^+[m-2] \psi_{N+m-1+i}^-[m-2] dx) \\
= & \delta_{i,j} C_{m-1+i}(t) + a_{i,j} - b a_{i,0} a_{0,j},
\end{aligned} \tag{4.6}$$

where

$$a_{i,j} = \int \psi_{m-1+j}^+[m-2] \psi_{N+m-1+i}^-[m-2] dx, \quad b = \frac{1}{C_{m-1}(t) + \int \psi_{m-1}^+[m-2] \psi_{N+m-1}^-[m-2] dx}.$$

Then

$$\begin{aligned}
& W_1(\psi_m^+[m-1], \dots, \psi_{m+k}^+[m-1]; \psi_{N+m}^-[m-1], \dots, \psi_{N+m+k}^-[m-1]; C_m(t), \dots, C_{m+k}(t)) = \\
& \begin{vmatrix} C_m(t) + a_{1,1} - b a_{1,0} a_{0,1} & a_{1,2} - b a_{1,0} a_{0,2} & \cdots & a_{1,k+1} - b a_{1,0} a_{0,k+1} \\ a_{2,1} - b a_{2,0} a_{0,1} & C_{m+1}(t) + a_{2,2} - b a_{2,0} a_{0,2} & \cdots & a_{2,k+1} - b a_{2,0} a_{0,k+1} \\ \vdots & \vdots & \ddots & \vdots \\ a_{k+1,1} - b a_{k+1,0} a_{0,1} & a_{k+1,2} - b a_{k+1,0} a_{0,2} & \cdots & C_{m+k}(t) + a_{k+1,k+1} - b a_{k+1,0} a_{0,k+1} \end{vmatrix}
\end{aligned}$$

$$\begin{aligned}
&= \begin{vmatrix} C_m(t) + a_{1,1} & a_{1,2} & \cdots & a_{1,k+1} \\ a_{2,1} & C_{m+1}(t) + a_{2,2} & \cdots & a_{2,k+1} \\ \vdots & \vdots & \ddots & \vdots \\ a_{k+1,1} & a_{k+1,2} & \cdots & C_{m+k}(t) + a_{k+1,k+1} \end{vmatrix} \\
&\quad - ba_{0,1} \begin{vmatrix} a_{1,0} & a_{1,2} & \cdots & a_{1,k+1} \\ a_{2,0} & C_{m+1}(t) + a_{2,2} & \cdots & a_{2,k+1} \\ \vdots & \vdots & \ddots & \vdots \\ a_{k+1,0} & a_{k+1,2} & \cdots & C_{m+k}(t) + a_{k+1,k+1} \end{vmatrix} \\
&\quad - ba_{0,2} \begin{vmatrix} C_m(t) + a_{1,1} & a_{1,0} & \cdots & a_{1,k+1} \\ a_{2,1} & a_{2,0} & \cdots & a_{2,k+1} \\ \vdots & \vdots & \ddots & \vdots \\ a_{k+1,1} & a_{k+1,0} & \cdots & C_{m+k}(t) + a_{k+1,k+1} \end{vmatrix} \\
&\quad - \cdots - ba_{0,k+1} \begin{vmatrix} C_m(t) + a_{1,1} & a_{1,2} & \cdots & a_{1,k} & a_{1,0} \\ a_{2,1} & C_{m+1}(t) + a_{2,2} & \cdots & a_{2,k} & a_{2,0} \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ a_{k+1,1} & a_{k+1,2} & \cdots & a_{k+1,k} & a_{k+1,0} \end{vmatrix} \\
&= b \begin{vmatrix} C_{m-1}(t) + a_{0,0} & a_{0,1} & a_{0,2} & \cdots & a_{0,k+1} \\ a_{1,0} & C_m(t) + a_{1,1} & a_{1,2} & \cdots & a_{1,k+1} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ a_{k+1,0} & a_{k+1,1} & a_{k+1,2} & \cdots & C_{m+k}(t) + a_{k+1,k+1} \end{vmatrix} =
\end{aligned}$$

$$\frac{W_1(\psi_{m-1}^+[m-2], \dots, \psi_{m+k}^+[m-2]; \psi_{N+m-1}^-[m-2], \dots, \psi_{N+m+k}^-[m-2]; C_{m-1}(t), \dots, C_{m+k}(t))}{C_{m-1}(t) + \int \psi_{m-1}^+[m-2] \psi_{N+m-1}^-[m-2] dx}.$$

So the formula (4.3a) holds. The formulas (4.3b) and (4.3c) can be proved similarly.

This completes the proof.

Theorem 4.1 Assume that $u, \Psi_1, \dots, \Psi_m, \Phi_1, \dots, \Phi_m$ is a solution of the KPESCS (3.5), $\psi_1^+, \dots, \psi_N^+$ and $\psi_{N+1}^-, \dots, \psi_{N+N}^-$ are solutions of (3.6) and (3.7) respectively, $C_1(t), \dots, C_N(t)$ are N arbitrary functions in t . Then the N -times repeated generalized binary Darboux transformation for (3.6) is given by

$$\psi^+[-N, +N] = \frac{W_2(\psi_1^+, \dots, \psi_N^+, \psi^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t))}{W_1(\psi_1^+, \dots, \psi_N^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t))}, \quad (4.7a)$$

$$u[-N, +N] = u + 2\partial^2 \ln W_1(\psi_1^+, \dots, \psi_N^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t)), \quad (4.7b)$$

$$\Psi_l[-N, +N] = -\frac{W_2(\psi_1^+, \dots, \psi_N^+, \Psi_l; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t))}{W_1(\psi_1^+, \dots, \psi_N^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t))}, \quad (4.7c)$$

$$\Phi_l[-N, +N] = -\frac{W_3(\psi_1^+, \dots, \psi_N^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-, \Phi_l; C_1(t), \dots, C_N(t))}{W_1(\psi_1^+, \dots, \psi_N^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t))}, \quad l = 1, \dots, m, \quad (4.7d)$$

$$\begin{aligned} \Psi_{m+j}[-N, +N] \\ = \frac{1}{2} \sqrt{\dot{C}_j(t)} \frac{W_2(\psi_1^+, \dots, \psi_{j-1}^+, \psi_{j+1}^+, \dots, \psi_N^+, \psi_j^+; \psi_{N+1}^-, \dots, \psi_{N+j-1}^-, \psi_{N+j+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_{j-1}(t), C_{j+1}(t), \dots, C_N(t))}{W_1(\psi_1^+, \dots, \psi_N^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t))}, \end{aligned} \quad (4.7e)$$

$$\begin{aligned} \Phi_{m+j}[-N, +N] \\ = \frac{1}{2} \sqrt{\dot{C}_j(t)} \frac{W_3(\psi_1^+, \dots, \psi_{j-1}^+, \psi_{j+1}^+, \dots, \psi_N^+, \psi_j^+; \psi_{N+1}^-, \dots, \psi_{N+j-1}^-, \psi_{N+j+1}^-, \dots, \psi_{N+N}^-, \psi_{N+j}^-; C_1(t), \dots, C_{j-1}(t), C_{j+1}(t), \dots, C_N(t))}{W_1(\psi_1^+, \dots, \psi_N^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t))}, \\ j = 1, \dots, N. \end{aligned} \quad (4.7f)$$

Namely,

$$\alpha \Psi_j^+[-N, +N]_y = -\Psi_j^+[-N, +N]_{xx} - u[-N, +N] \Psi_j^+[-N, +N], \quad (4.8a)$$

$$\alpha \Phi_j^-[-N, +N]_y = \Phi_j^-[-N, +N]_{xx} + u[-N, +N] \Phi_j^-[-N, +N], \quad j = 1, \dots, m + N, \quad (4.8b)$$

$$\alpha \psi^+[-N, +N]_y = -\psi^+[-N, +N]_{xx} - u[-N, +N] \psi^+[-N, +N], \quad (4.8c)$$

$$\psi^+[-N, +N]_t = A^+(u[-N, +N]) \psi^+[-N, +N] + T_{m+N}^+(\Psi[-N, +N], \Phi[-N, +N]) \psi^+[-N, +N]. \quad (4.8d)$$

So $u[-N, +N]$, $\Psi_j[-N, +N]$, $\Phi_j[-N, +N]$, $j = 1, \dots, m + N$ satisfy the KPESCS (3.5) with degree $(m + N)$.

Proof: For notation simplicity, we denote $\psi^+[-l, +l]$, $u[-l, +l]$, $\Psi_i[-l, +l]$, $\Phi_i[-l, +l]$ by $\psi^+[l]$, $u[l]$, $\Psi_i[l]$, $\Phi_i[l]$ respectively, $0 \leq l \leq N$, $1 \leq i \leq m + N$. By (3.21) and (4.3),

$$\begin{aligned} \psi^+[N] &= \psi^+[N-1] - \frac{\psi_N^+[N-1] \int \psi_{N+N}^-[N-1] \psi^+[N-1] dx}{C_N(t) + \int \psi_{N+N}^-[N-1] \psi_N^+[N-1] dx} \\ &= \frac{W_2(\psi_N^+[N-1], \psi^+[N-1]; \psi_{N+N}^-[N-1]; C_N(t))}{W_1(\psi_N^+[N-1]; \psi_{N+N}^-[N-1]; C_N(t))} \\ &= \frac{W_2(\psi_{N-1}^+[N-2], \psi_N^+[N-2], \psi^+[N-2]; \psi_{N+N-1}^-[N-2], \psi_{N+N}^-[N-2]; C_{N-1}(t), C_N(t))}{C_{N-1}(t) + \int \psi_{N+N-1}^-[N-2] \psi_{N-1}^+[N-2] dx} \\ &\quad \times \frac{C_{N-1}(t) + \int \psi_{N+N-1}^-[N-2] \psi_{N-1}^+[N-2] dx}{W_1(\psi_{N-1}^+[N-2], \psi_N^+[N-2]; \psi_{N+N-1}^-[N-2], \psi_{N+N}^-[N-2]; C_{N-1}(t), C_N(t))} \\ &= \dots \\ &= \frac{W_2(\psi_1^+, \dots, \psi_N^+, \psi^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t))}{W_1(\psi_1^+, \dots, \psi_N^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t))}, \end{aligned}$$

so (4.7a) holds, Similarly we can prove (4.7c) and (4.7d) hold,

$$\begin{aligned}
u[N] &= u[N-1] + 2\partial^2 \ln(C_N(t) + \int \psi_N^+[N-1] \psi_{N+N}^-[N-1] dx) \\
&= u[N-1] + 2\partial^2 \ln W_1(\psi_N^+[N-1]; \psi_{N+N}^-[N-1]; C_N(t)) \\
&= u[N-2] + 2\partial^2 \ln(C_{N-1}(t) + \int \psi_{N-1}^+[N-2] \psi_{N+N-1}^-[N-2] dx) + \\
&\quad 2\partial^2 \ln \left[\frac{W_1(\psi_{N-1}^+[N-2], \psi_N^+[N-2]; \psi_{N+N-1}^-[N-2], \psi_{N+N}^-[N-2]; C_{N-1}(t), C_N(t))}{C_{N-1}(t) + \int \psi_{N-1}^+[N-2] \psi_{N+N-1}^-[N-2] dx} \right] \\
&= u[N-2] + 2\partial^2 \ln W_1(\psi_{N-1}^+[N-2], \psi_N^+[N-2]; \psi_{N+N-1}^-[N-2], \psi_{N+N}^-[N-2]; \\
&\quad C_{N-1}(t), C_N(t)) \\
&= \dots \\
&= u + 2\partial^2 \ln W_1(\psi_1^+, \dots, \psi_N^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t)),
\end{aligned}$$

so (4.7b) holds.

From (3.21a) and (3.21e), we have

$$\psi_j^+[j] = \psi_j^+[j-1] - \frac{\psi_j^+[j-1] \int \psi_j^+[j-1] \psi_{N+j}^-[j-1] dx}{C_j(t) + \int \psi_j^+[j-1] \psi_{N+j}^-[j-1] dx} = \frac{C_j(t) \psi_j^+[j-1]}{C_j(t) + \int \psi_j^+[j-1] \psi_{N+j}^-[j-1] dx},$$

$$\Psi_{m+j}[j] = \frac{1}{2} \frac{\sqrt{\dot{C}_j(t)} \psi_j^+[j-1]}{C_j(t) + \int \psi_j^+[j-1] \psi_{N+j}^-[j-1] dx},$$

so

$$\Psi_{m+j}[j] = \frac{1}{2} \frac{\sqrt{\dot{C}_j(t)}}{C_j(t)} \psi_j^+[j].$$

Analogously as above, we have

$$\begin{aligned}
&\Psi_{m+j}[N] \\
&= \frac{W_2(\psi_N^+[N-1], \Psi_{m+j}[N-1]; \psi_{N+N}^-[N-1]; C_N(t))}{W_1(\psi_N^+[N-1]; \psi_{N+N}^-[N-1]; C_N(t))} \\
&= \frac{W_2(\psi_{j+1}^+[j], \dots, \psi_N^+[j], \Psi_{m+j}[j]; \psi_{N+j+1}^-[j], \dots, \psi_{N+N}^-[j]; C_{j+1}(t), \dots, C_N(t))}{W_1(\psi_{j+1}^+[j], \dots, \psi_N^+[j]; \psi_{N+j+1}^-[j], \dots, \psi_{N+N}^-[j]; C_{j+1}(t), \dots, C_N(t))} \\
&= \frac{\sqrt{\dot{C}_j(t)}}{2C_j(t)} \frac{W_2(\psi_{j+1}^+[j], \dots, \psi_N^+[j], \psi_j^+[j]; \psi_{N+j+1}^-[j], \dots, \psi_{N+N}^-[j]; C_{j+1}(t), \dots, C_N(t))}{W_1(\psi_{j+1}^+[j], \dots, \psi_N^+[j]; \psi_{N+j+1}^-[j], \dots, \psi_{N+N}^-[j]; C_{j+1}(t), \dots, C_N(t))} \\
&= \frac{\sqrt{\dot{C}_j(t)}}{2C_j(t)} \frac{W_2(\psi_1^+, \dots, \psi_N^+, \psi_j^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t))}{W_1(\psi_1^+, \dots, \psi_N^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t))} \\
&= \frac{1}{2} \sqrt{\dot{C}_j(t)} \frac{W_2(\psi_1^+, \dots, \psi_{j-1}^+, \psi_{j+1}^+, \dots, \psi_N^+, \psi_j^+; \psi_{N+1}^-, \dots, \psi_{N+j-1}^-, \psi_{N+j+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_{j-1}(t), C_{j+1}(t), \dots, C_N(t))}{W_1(\psi_1^+, \dots, \psi_N^+; \psi_{N+1}^-, \dots, \psi_{N+N}^-; C_1(t), \dots, C_N(t))},
\end{aligned}$$

so (4.7e) holds. In a similar way we can prove (4.7f) holds. This completes the proof.

Remark:

(1). If $C_1(t), \dots, C_N(t)$ are replaced by N arbitrary constants C_1, \dots, C_N , then (4.7) gives the N -times repeated Darboux transformation for (3.18).

(2). The N -times repeated generalized binary Darboux transformation (4.7) provides a non-auto Bäcklund transformation between two KPESCSs with degrees m and $(m + N)$ respectively. We now use this to construct some interesting solutions for the KPESCS.

Example 4: (Degenerate case)

When some $C_j(t)$ in (4.7) are chosen to be arbitrary numerical constants C_j , we can get the some interesting solutions for the KPESCS (3.5) in the degenerate case ("degenerate" means the number of soliton is larger than the degree of the equation). In the following, we will take the KPIIESCS (3.32) for example and choose $u = 0$ as the initial solution of it with $m = 0$.

Example 4.1: two-soliton solution of the KPIIESCS.

Take

$$\begin{aligned}\psi_i^+ &= e^{\eta_i}, \quad \psi_{2+i}^- = e^{\theta_i}, \quad i = 1, 2, \\ \eta_i &= l_i x - l_i^2 y - 4l_i^3 t, \quad \theta_i = k_i x + k_i^2 y - 4k_i^3 t, \quad l_i, k_i \in \mathbb{R}, \\ C_1(t) &= \frac{1}{k_1 + l_1} e^{2\beta(t)}, \quad C_2(t) \equiv \frac{c}{k_2 + l_2},\end{aligned}$$

where $\beta(t)$ is an arbitrary function in t and c is an arbitrary constant. From (4.7b), (4.7e) and (4.7f), we get the 2-soliton solution of the KPIIESCS (3.32) with $m = 1$ as follows

$$\begin{aligned}& u(x, y, t)[-2, +2] \\ &= 2\partial^2 \ln W_1(\psi_1^+, \psi_2^+; \psi_3^-, \psi_4^-; C_1(t), C_2(t)) \\ &= 2\partial^2 \ln \left| \begin{array}{cc} C_1(t) + \int \psi_3^- \psi_1^+ dx & \int \psi_3^- \psi_2^+ dx \\ \int \psi_4^- \psi_1^+ dx & C_2(t) + \int \psi_4^- \psi_2^+ dx \end{array} \right| \\ &= 2\partial^2 \ln \left| \begin{array}{cc} \frac{1}{k_1 + l_1} e^{2\beta(t)} + \frac{1}{k_1 + l_1} e^{\theta_1 + \eta_1} & \frac{1}{k_1 + l_2} e^{\theta_1 + \eta_2} \\ \frac{1}{k_2 + l_1} e^{\theta_2 + \eta_1} & \frac{c}{k_2 + l_2} e^{\theta_2 + \eta_2} \end{array} \right| \\ &= 2\partial^2 \ln \left[\frac{ce^{2\beta(t)}}{(k_1 + l_1)(k_2 + l_2)} (1 + e^{\theta_1 + \eta_1 - 2\beta(t)} + e^{\theta_2 + \eta_2 + \xi_0} + \frac{(k_1 - k_2)(l_1 - l_2)}{(k_2 + l_1)(k_1 + l_2)} e^{\theta_1 + \theta_2 + \eta_1 + \eta_2 - 2\beta(t) + \xi_0}) \right] \\ &= 2\partial^2 \ln [1 + e^{\theta_1 + \eta_1 - 2\beta(t)} + e^{\theta_2 + \eta_2 + \xi_0} + \frac{(k_1 - k_2)(l_1 - l_2)}{(k_2 + l_1)(k_1 + l_2)} e^{\theta_1 + \theta_2 + \eta_1 + \eta_2 - 2\beta(t) + \xi_0}],\end{aligned}$$

where $\xi_0 = -\ln c$,

$$\begin{aligned}& \Psi_1[-2, +2] \\ &= \frac{1}{2} \sqrt{\dot{C}_1(t)} \frac{W_2(\psi_2^+, \psi_1^+; \psi_4^-, \psi_3^-; C_2(t))}{W_1(\psi_1^+, \psi_2^+; \psi_3^-, \psi_4^-; C_1(t), C_2(t))} \\ &= \frac{1}{2} \sqrt{2(k_1 + l_1) \dot{\beta}(t)} e^{\eta_1 - \beta(t)} \frac{1 + \frac{l_1 - l_2}{k_2 + l_1} e^{\theta_2 + \eta_2 + \xi_0}}{1 + e^{\theta_1 + \eta_1 - 2\beta(t)} + e^{\theta_2 + \eta_2 + \xi_0} + \frac{(k_1 - k_2)(l_1 - l_2)}{(k_2 + l_1)(k_1 + l_2)} e^{\theta_1 + \theta_2 + \eta_1 + \eta_2 - 2\beta(t) + \xi_0}}, \\ & \Phi_1[-2, +2] \\ &= \frac{1}{2} \sqrt{\dot{C}_1(t)} \frac{W_3(\psi_2^+; \psi_4^-, \psi_3^-; C_2(t))}{W_1(\psi_1^+, \psi_2^+; \psi_3^-, \psi_4^-; C_1(t), C_2(t))} \\ &= \frac{1}{2} \sqrt{2(k_1 + l_1) \dot{\beta}(t)} e^{\theta_1 - \beta(t)} \frac{1 + \frac{k_1 - k_2}{k_1 + l_2} e^{\theta_2 + \eta_2 + \xi_0}}{1 + e^{\theta_1 + \eta_1 - 2\beta(t)} + e^{\theta_2 + \eta_2 + \xi_0} + \frac{(k_1 - k_2)(l_1 - l_2)}{(k_2 + l_1)(k_1 + l_2)} e^{\theta_1 + \theta_2 + \eta_1 + \eta_2 - 2\beta(t) + \xi_0}},\end{aligned}$$

$$\Psi_2[-2, +2] = \Phi_2[-1, +2] = 0.$$

In a similar way, for $\forall N, m \in \mathbb{N}$, $N > m$, when $C_j(t), j = 1, \dots, m$ are taken to be arbitrary functions in t and $C_j(t), j = m+1, \dots, N$ are taken to be numerical constants, we can get the N -soliton solution of (3.32) with degree m .

Example 4.2: mixture of the exponential and rational solutions of the KPIESCS.

When take

$$\psi_1^+ = e^\zeta(x - 2qy - 12q^2t), \quad \psi_2^+ = e^\zeta, \quad \zeta = qx - q^2y - 4q^3t, \quad q \in \mathbb{R},$$

$$\psi_3^- = e^\xi(x + 2ky - 12k^2t), \quad \psi_4^- = e^\xi, \quad \xi = kx + k^2y - 4k^3t, \quad k \in \mathbb{R},$$

$$C_1(t) = \frac{1}{q+k}e^{2\beta(t)}, \quad C_2(t) \equiv \frac{c}{q+k},$$

where $\beta(t)$ is an arbitrary function in t and c is a constant, we will get another solution (mixture of the exponential and rational solutions) of the KPIESCS with degree $m = 1$ as follows

$$u(x, y, t)[-2, +2] = 2\partial^2 \ln[1 + e^{\zeta+\xi+\xi_0} + e^{\zeta+\xi-2\beta(t)}r + \frac{1}{(k+q)^2}e^{2(\zeta+\xi)-2\beta(t)+\xi_0}],$$

$$\Psi_1[-2, +2] = \frac{1}{2}\sqrt{2(k+q)\dot{\beta}(t)}e^{\zeta-\beta(t)}\frac{x - 2qy - 12q^2t + \frac{1}{k+q}e^{\zeta+\xi+\xi_0}}{1 + e^{\zeta+\xi+\xi_0} + e^{\zeta+\xi-2\beta(t)}r + \frac{1}{(k+q)^2}e^{2(\zeta+\xi)-2\beta(t)+\xi_0}},$$

$$\Phi_1[-2, +2] = \frac{1}{2}\sqrt{2(k+q)\dot{\beta}(t)}e^{\xi-\beta(t)}\frac{x + 2ky - 12k^2t + \frac{1}{k+q}e^{\zeta+\xi+\xi_0}}{1 + e^{\zeta+\xi+\xi_0} + e^{\zeta+\xi-2\beta(t)}r + \frac{1}{(k+q)^2}e^{2(\zeta+\xi)-2\beta(t)+\xi_0}},$$

where $\xi_0 = -\ln c$, $r = (x + 2ky - 12k^2t - \frac{1}{k+q})(x - 2qy - 12q^2t - \frac{1}{k+q}) + \frac{1}{(k+q)^2}$.

Example 5: (Nondegenerate case) N -soliton solution.

Take the case of $\alpha = i$ for example. We take $u = 0$ as the initial solution of (3.26) with $m = 0$ and let

$$\psi_j^+ = e^{k_j x + i k_j^2 y - 4k_j^3 t}, \quad \psi_{N+j}^- = \overline{\psi_j^+}, \quad k_j \in \mathbb{C},$$

$$C_j(t) = \frac{1}{2\operatorname{Re}(k_j)}e^{2\beta_j(t)}, \quad j = 1, \dots, N,$$

where $\beta_j(t)$ are arbitrary functions in t . Then the formulas (4.7b), (4.7e) and (4.7f) give the N -soliton solution of the KPIESCS (3.26) with degree N .

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References

- [1] Mel'nikov V K 1988 Phys. Lett. A 133 493
- [2] Mel'nikov V K 1989 Commun. Math. Phys. 120 451
- [3] Mel'nikov V K 1990 J. Math. Phys. 31 1106
- [4] Leon J, Latifi A 1990 J. Phys. A 23 1385-1403
- [5] Leon J 1990 Phys. Lett. A 144 444-452
- [6] Doktorov E V, Vlasov R A 1991 Opt. Acta 30 3321
- [7] Mel'nikov V K 1992 Inverse Probl. 8 133
- [8] Zeng Y B , Ma W X and Lin R L 2000 J. Math. Phys. 41 (8) 5453-5489
- [9] Lin R L, Zeng Y B and Ma W X 2001 Physica A 291 287-298
- [10] Ye S, Zeng Y B 2002 J. Phys. A 35 L283-L291
- [11] Zeng Y B, Ma W X and Shao Y J 2001 J. Math. Phys. 42(5) 2113-2128
- [12] Zeng Y B, Shao Y J and Ma W X 2002 Commun. Theor. Phys. 38 641-648
- [13] Zeng Y B, Shao Y J and Xue W M 2003 J. Phys. A 36 1-9
- [14] Mel'nikov V K 1987 Commun. Math. Phys. 112 639-652
- [15] Mel'nikov V K 1989 Commun. Math. Phys. 126 201-215
- [16] Deng S F, Chen D Y and Zhang D J 2003 J. Phys. Soc. Jap. 72 2184-2192
- [17] Ablowitz M J, Clarkson P A 1991 Solitons, Nonlinear Evolution Equations and Inverse Scattering (Cambridge)
- [18] Dickey L A 1991 Soliton equation and Hamiltonian systems (Singapore:World Scientific)
- [19] Matveev V B, Salle M A 1991 Darboux Transformations and Solitons (Berlin: Springer)

- [20] Date E, Jimbo M, Kashiwara M and Miwa T 1983 In Nonlinear Integrable Systems-
Classical Theory and Quantum Theory, Jimbo M, Miwa T 1983 (eds.) (Singapore:World Scientific)
- [21] Ohta Y, Satsuma J, Takahashi D and Tokihiro T 1988 Prog. Theor. Phys. Suppl. 94
210
- [22] Jurij Sidorenko, Walter Strampp 1991 Inverse Probl. 7 L37-L43
- [23] Oevel W, Strampp W 1993 Commun. Math. Phys. 157 51-81
- [24] Dickey L A 1995 Lett. Math. Phys. 34 379-384
- [25] Cheng Y 1992 J. Math. Phys. 33 3774
- [26] Cheng Y 1995 Commun. Math. Phys. 171 661-682